

Computational statistical physics: The example of hard spheres

Lectures at the 16th Chris Engelbrecht Summer School in
Theoretical Physics,
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Fourth part: Entropic phase transitions

Werner Krauth

Laboratoire de Physique Statistique
Ecole Normale Supérieure, Paris, France

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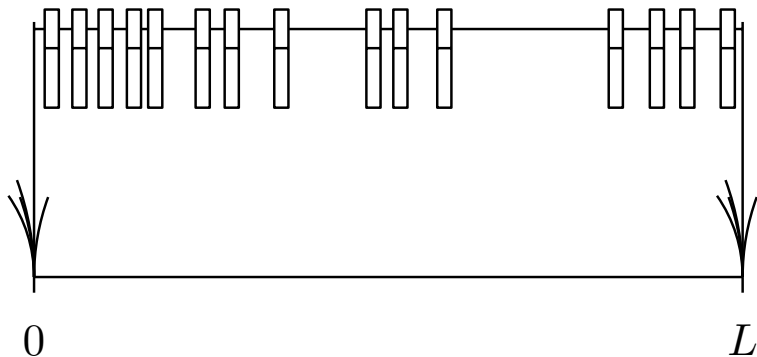




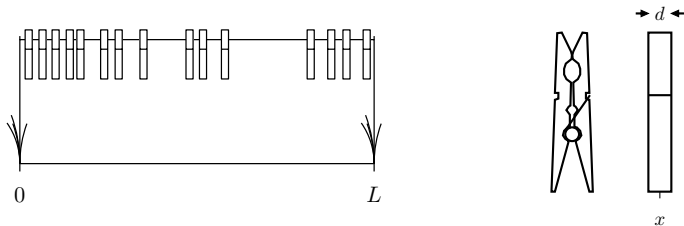
- Origin of ‘forces’ in systems without interactions



Random clothes-pin model



Random clothes-pin model



- N pins, width $d = 2\sigma$, length L .
- simplest entropic system (\equiv one-dimensional hard disks)

$$\pi(x_1, \dots, x_N) = \begin{cases} 1 & \text{if no overlap} \\ 0 & \text{otherwise} \end{cases}$$

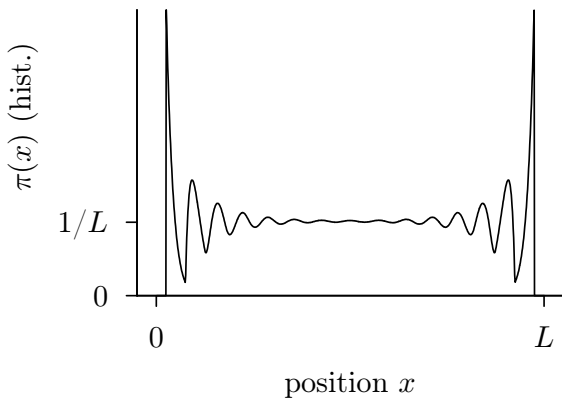


```
procedure naive-pins
1 for  $k = 1, \dots, N$  do
   {  $x_k \leftarrow \text{ran}[d/2, L - d/2]$ 
     for  $l = 1, \dots, k - 1$  do
       { if  $(|x_k - x_l| < d)$  goto 1 (reject sample—tabula rasa)
     }
   output  $\{x_1, \dots, x_N\}$ 
```

- What is the density $\pi(x)$ of clothes-pins for $0 < x < L$?
- NB: $\pi(x)$: probability to have a pin-center at x .



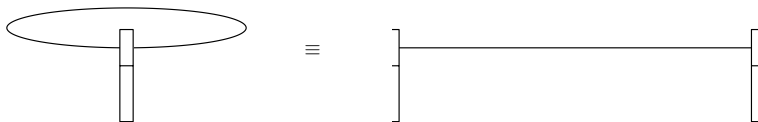
Output of Alg.naive-pins



- Just a boundary effect?



Not a boundary effect

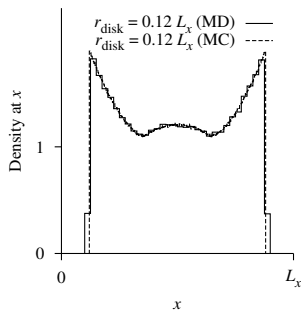
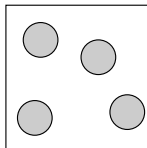


$$\underbrace{\pi(x_k, x_l) = \pi_2(x_k - x_l)}_{\text{pair correlation periodic system}} = \underbrace{\pi(x_k - d/2)}_{\text{density segment}}$$

- density of periodic system is constant (translation symmetry)



Boundary effects in two dimensions



Clothes-pin partition function

$$\begin{aligned} Z_{\text{pin}}(N, L) &= \int_0^L dx_1 \dots \int_0^L dx_N \overbrace{\pi(x_1, \dots, x_N)}^{\text{totally symmetric in } x_1, \dots, x_N} \\ &= N! \int_0^L dx_1 \dots \int_0^L dx_N \underbrace{\pi(x_1, \dots, x_N)}_{x_1 < x_2 < \dots < x_N} \\ &= N! \int_{d/2}^{L+\frac{d}{2}-Nd} dx_1 \int_{x_1+d}^{L+\frac{d}{2}-(N-1)d} dx_2 \dots \int_{x_{N-1}+d}^{L-d/2} dx_N \end{aligned}$$

transformation $y_k = x_k - x_{k-1} - d$ with $\partial y_k / \partial x_k = 1$

$$\begin{aligned} &= N! \int_0^{L-Nd} dy_1 \dots \int_0^{L-Nd} dy_N \{y_1 < \dots < y_N\} \\ &= \dots \end{aligned}$$



Clothes-pin partition function

$$\begin{aligned} Z_{\text{pin}}(N, L) &= \\ &= N! \int_0^{L-Nd} dy_1 \dots \int_0^{L-Nd} dy_N \{y_1 < \dots < y_N\} \\ &= \dots \end{aligned}$$



Clothes-pin partition function II

$$Z(N, L) = \begin{cases} (L - d \cdot N)^N & \text{if } L > d \cdot N \\ 0 & \text{otherwise} \end{cases}$$

...

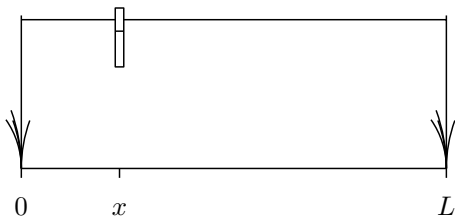


```
procedure direct-pin  
 $\Delta \leftarrow L - d \cdot N$   
for  $k = 1, \dots, N$  do  
  {  $\tilde{y}_k \leftarrow \text{ran}[0, \Delta]$   
  {  $y_1, \dots, y_N \leftarrow \text{sort}[\{\tilde{y}_1, \dots, \tilde{y}_N\}]$   
  for  $k = 1, \dots, N$  do  
    {  $x_k \leftarrow y_k + d/2 + d \cdot (k - 1)$   
  output {  $x_1, \dots, x_N$  }
```

- works at **any** density, without rejections, unlike Alg.direct-disks



computing $\pi(x)$ from partition function $Z(k, x)$

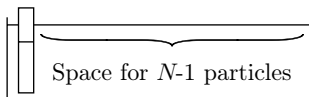
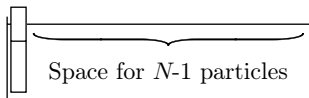


$$\begin{aligned}\pi(x) &= \sum_{k=0}^{N-1} \pi_k(x) = \\ &= \sum_{k=0}^N \frac{1}{Z(N, L)} \binom{N-1}{k} Z(k, x - \frac{d}{2}) Z(N-1-k, L-x - \frac{d}{2})\end{aligned}$$

- $\pi(x)$ has a **strictly** constant piece, if $N \cdot d < L/2!!$



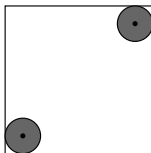
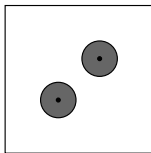
Depletion interaction



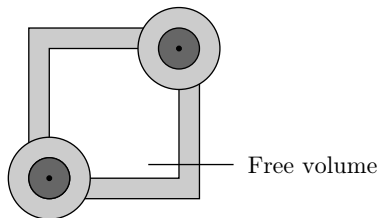
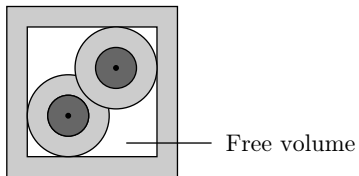
- attractive **and** repulsive
- two-body interaction only in one dimension



Analysis of spatial repartition



Entropic Attraction



Depletion interaction

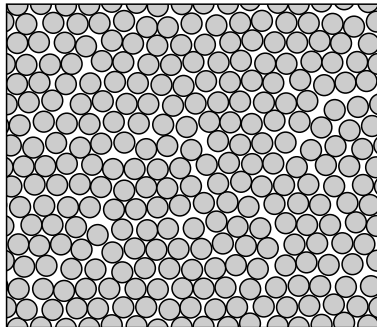
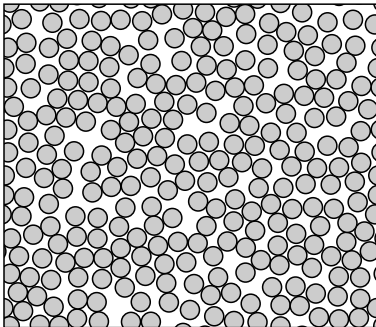
- Asakura-Oosawa 'depletion' interaction
- of fundamental importance in soft condensed matter (polymers, colloids etc)
- underlies liquid-solid phase transition in hard spheres
- ... and phase-separation transition in binary mixtures

The hard sphere paradox

Hard spheres are not impenetrable ?!?!?!?



Large systems of hard spheres



- Liquid–Solid phase transition in dimensions 2 or higher.



Binary mixtures

